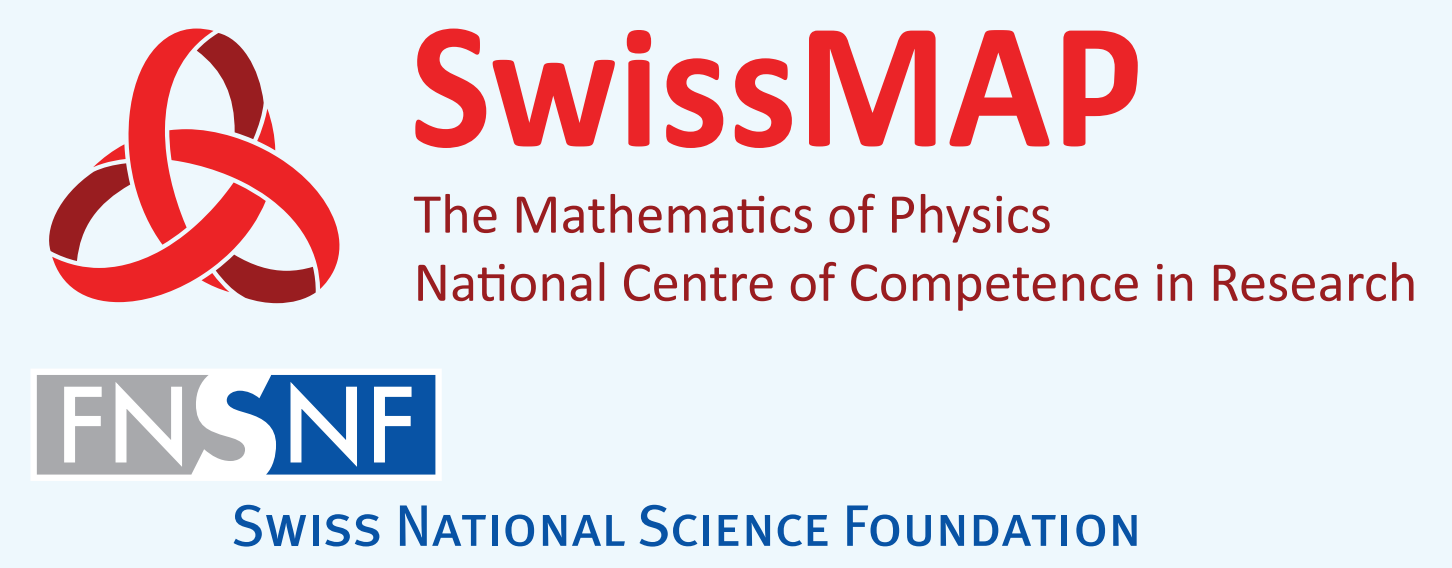




EIGENSTATE THERMALISATION & ERGODICITY IN LOW DIMENSIONAL THEORIES

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1. BACKGROUND & MOTIVATION

- The Eigenstate Thermalization Hypothesis (ETH) provides a possible way to explain thermalization in a quantum system.
- It states that expectation value of a 'typical' operator in an Eigenstate is approximately same as its expectation value in a micro-canonical ensemble, up to exponentially suppressed corrections in system size/entropy.
- Random Matrix Theory (RMT) is a canonical example of a system obeying ETH.
- Wigner-Dyson (WD) spectral statistics is a notable attribute of RMT. These spectral statistics are also underlying property of Quantum Chaotic systems.
- The SYK model, a model of large N number of Majorana Fermions, provides a playground to test ETH analytically. It has a universal low energy limit described by 'Schwarzian theory'.
- We analytically demonstrate ETH in the universal Schwarzian theory, as well as in the SYK in various limits.
- Away from the ergodic regime, where WD statistics hold, the corrections are non-universal.
- We study the behaviour of the operators in the ergodic regime of the SYK model.
- Using the AdS/CFT correspondence we can relate the black hole formation in a theory of gravity to thermalisation in the dual field theory.
- Our ultimate goal is to develop a microscopic understanding of black hole formation by improving our understanding of thermalisation.

2. INTRODUCTION TO THE SYK MODEL

1. SYK Hamiltonian [1]:

$$H = (i)^{\frac{q}{2}} \sum_{1 \leq i_1 < i_2 \dots < i_q \leq N} j_{i_1 \dots i_q} \psi_{i_1} \psi_{i_2} \dots \psi_{i_q}, \quad \langle j_{i_1 \dots i_q}^2 \rangle = \frac{J^2 (q-1)!}{N^{q-1}}. \quad (1)$$

2. **IR Effective action:** In the IR limit, $J|\tau| \rightarrow \infty$, reparametrization is an emergent symmetry of the action. This is spontaneously broken by the ground state, and explicitly broken by introducing large but finite J . The low-energy effective action is given by Schwarzian action:

$$S_{Sch} \sim \frac{N}{J} \int d\tau \{f(\tau), \tau\}, \quad \{f(\tau), \tau\} = \frac{f'''(\tau)}{f'(\tau)} - \frac{3}{2} \left(\frac{f''(\tau)}{f'(\tau)} \right)^2 \quad (2)$$

3. **Computation of the correlation functions:** In the large N limit, correlations functions of fermions receive contributions from,

- Sum over ladders
- Sum over $f(\tau)$ modes

3. ETH IN THE CONFORMAL LIMIT

1. In the IR limit, the conformal four-point function can be computed exactly by summing over the ladder diagrams.

2. Spectrum of operators in this limit is:

$$\mathcal{O}_n = \frac{1}{\sqrt{N}} \sum_{k=0}^{2n+1} \left[\sum_{i=1}^N d_{nk} \partial^k \psi_i \partial^{2n+1-k} \psi_i \right], \quad (3)$$

with conformal weights, $h_n = 2n + 1 + 2\epsilon_n$

3. Correlation functions of operators \mathcal{O}_n were computed in [2] for arbitrary values of m, n, k and q .

4. We study the three-point function, $\langle \mathcal{O}_m \mathcal{O}_k \mathcal{O}_n \rangle$ in the limit $m, n \gg k \sim \mathcal{O}(1)$, (where $E = \frac{m+n}{2}$, $d = n - m$), [5],

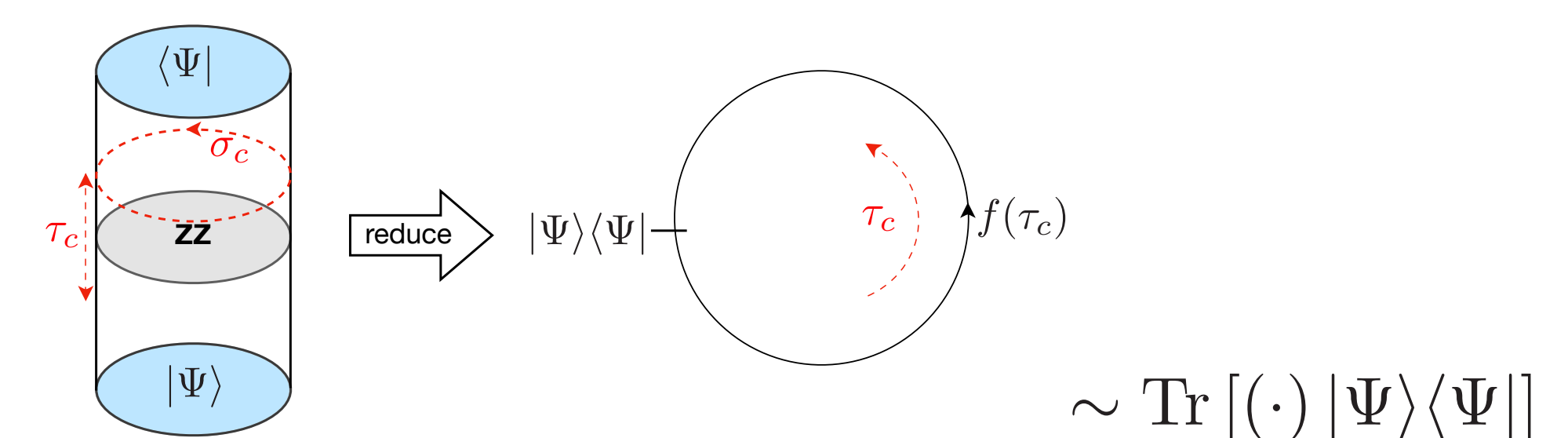
$$\langle \mathcal{O}_m \mathcal{O}_k \mathcal{O}_n \rangle = \mathfrak{g}(E, d) \times \left[2^{-2(2E-1)} \sqrt{2\pi E} \frac{\Gamma(4E-1)}{\Gamma(2E-d)\Gamma(2E+d)} \right] \quad (4)$$

$$E \rightarrow \infty \quad \boxed{f_k(E) \delta_{m,n} + e^{-E \ln^2} R(E, d)}$$

This is the same behaviour for ETH as discussed in [3]

4. ETH IN THE SCHWARZIAN LIMIT

1. The Schwarzian action, (2), can be obtained from a dimensional reduction from 2-dimensional boundary Liouville theory, [4, 6].



2. Depending on boundary conditions of the boundary Liouville CFT, we obtain an operator/density matrix insertion in the Schwarzian path integral. These correspond to pure coherent states of the Schwarzian theory, [6].

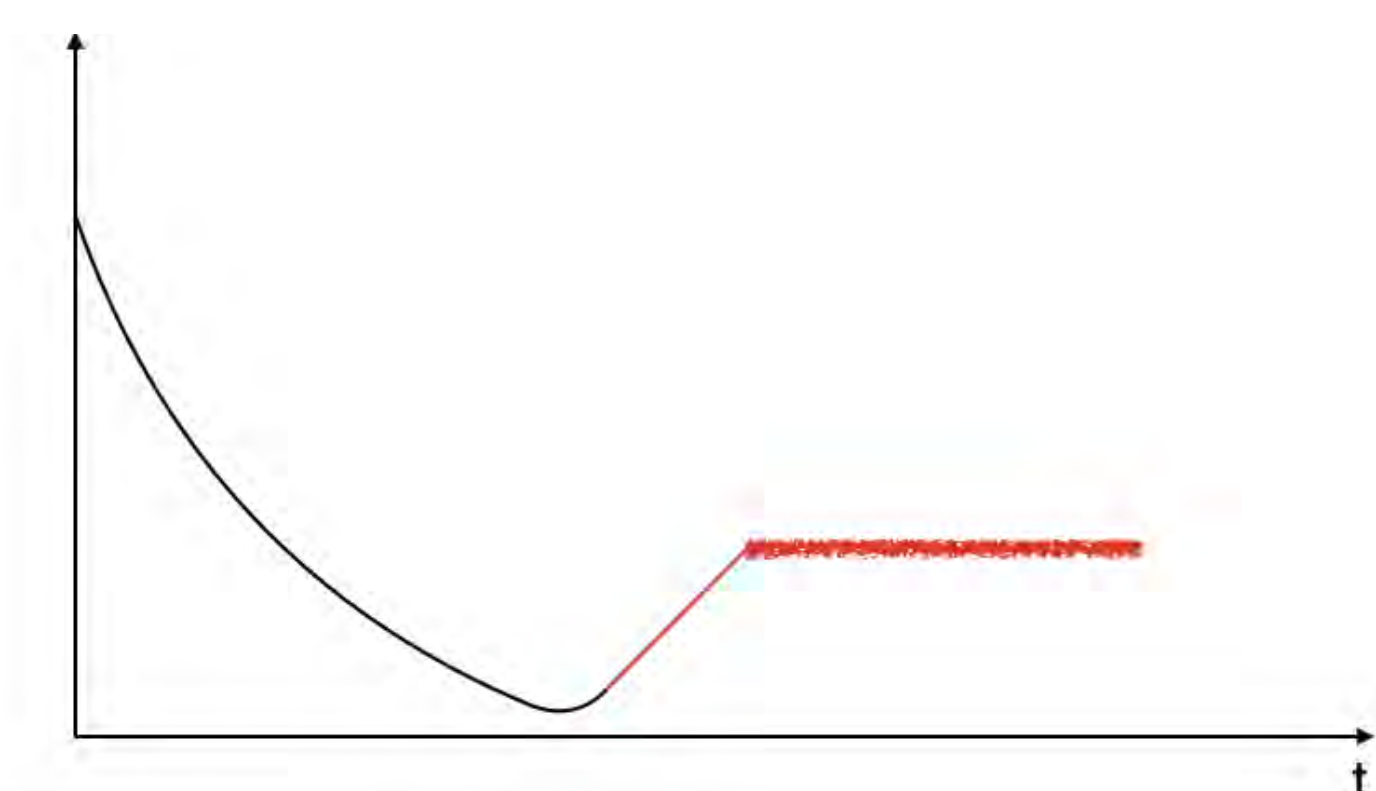
3. Summary of thermal nature of different states:

| state | 2D picture | class | ETH | λ |
|---------------------------------|------------|------------|-----|-------------------|
| $ E(k)\rangle$ | ZZ | parabolic | ✓ | $2\pi T_{ETH}$ |
| $ E_{r-}\rangle$ | FZZT | elliptic | ✓ | $2\pi T_{ETH}$ |
| $ E_{r+}\rangle$ | FZZT | hyperbolic | ✗ | $\in i\mathbb{R}$ |
| $\mathcal{O}_{\ell_H} 0\rangle$ | ZZ | parabolic | ✓ | — |

5. ERGODICITY IN THE SYK MODEL

1. Using the 1st quantized Hamiltonian of the SYK model, the ergodic limit of the SYK model can be found, [7]. In this limit, one can analytically extract WD from the SYK Hamiltonian.

2. Operator correlation functions have a characteristic behaviour in this regime,



3. In an ongoing work, we study the ergodic limit of the operator correlation functions in the SYK model.

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